Semiconductor scintillator based on photon recycling

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1. Introduction

The key issue in implementing a semiconductor scintillator is to make sure that photons generated deep inside the semiconductor slab could reach its surface without tangible attenuation. However, semiconductors are usually opaque at wavelengths corresponding to their radiative emission spectrum. Our group has been working on the implementation of scintillators based on direct-gap semiconductors, like InP or GaAs. For the exemplary case of InP the luminescence spectrum is a band of wavelengths near 920 nm.

The original idea [1] was to make InP relatively transparent to this radiation by doping it heavily with donor impurities, so as to introduce the Burstein shift between the emission and the absorption spectra. The problem with this approach is attenuation of the optical signal by the free-carrier absorption (FCA) in heavily-doped material.

Here we shall describe another approach, based on the extremely high quantum radiative efficiency of high-quality direct-gap semiconductors, such as InP. In these materials, an act of interband absorption does not finish off the luminescence but merely creates a new minority carrier, which in turn recombines in a predominately radiative fashion. To take full advantage of this “photon recycling” effect, optically-tight integration of photoreceivers on both sides of the semiconductor scintillator wafer is desirable.

2. Photon recycling

Consider an InP scintillator slab with two photoreceiver systems integrated on the opposite sides of the slab [7]. Let the interaction of strength G occur a distance z from the detector top surface, as indicated in Fig. 1. The minority carrier (“hole”) h has the probability \( P(z) \) to be on par with that of diode detectors implemented in the same material.

The proportionality of scintillation yield is not the only expected advantage of semiconductor scintillators. One of the major benefits of semiconductor materials is the mature technology that enables the implementation of epitaxial photodiodes integrated on the surface of a semiconductor slab [5]. The epitaxial diode provides nearly perfect registration efficiency of photons that have reached the heterointerface.

A semiconductor scintillator endowed with an integrated photoreceiver can be patterned into a two-dimensional array of pixels. Such an array forms a basic unit that can be stacked up indefinitely in the third direction. This enables three-dimensional integration of scintillator “voxels” and thus offers a tantalizing possibility of implementing a compact low-voltage Compton telescope [6].
The detector signals $D_1$ and $D_2$ add single-pass contributions from different cycles. As is evident from Fig. 1, the sum can be found as geometric progression, giving

$$D_i(z) = Gmp_i(z) \times \sum_{n=0}^{\infty} \left[\eta(1-P)\right]^n = \frac{Gmp_i(z)}{\xi + \eta P(z)} \quad i = 1, 2$$

and the total photon collection efficiency, $\text{PCE} = (D_1 + D_2)/G$, is given by

$$\text{PCE} = \frac{p_1(z) + p_2(z)}{[1/(\eta + \beta)] + [p_1(z) + p_2(z)]}$$

We note that for high photon recycling ($\eta \approx 1$ and $\beta \approx 0$) one has an ideal scintillator in the sense that the entire generated luminescence is collected – even though the single-pass probabilities $p_1$ and $p_2$ may not be high due to interband absorption.

Fig. 2 shows the photon collection efficiency for InP scintillator doped $n$-type with the concentration $N_n = 3 \times 10^{17} \text{ cm}^{-3}$. The calculation is based on Eqs. (4) and (5) and measured quantum efficiency $\eta$ (Fig. 3a) and the FCA absorption coefficient $\chi_a$ (Fig. 3b).

The only approximation involved in Eq. (4) is the assumption that every act of recycling occurs at the same place $z$ where the initial interaction occurred, and therefore the same probabilities $p_1(z)$ and $p_2(z)$ appear at all stages of the recycling, see Fig. 1. This has reduced the summation of an infinite series to a geometric progression and allowed us to obtain the result in a closed form. In reality, however, there is photon-assisted transport of holes in photon recycling, which has the nature of a random walk [6]. We have evaluated this effect (to be published separately) and found that its inclusion does not change the results qualitatively, although it does slightly enhance the estimate of PCE.

As seen from Fig. 2a, photon recycling delivers a reasonable fraction of the scintillating photons to the wafer surface. However, this fraction depends on the exact position of the interaction.

Let the luminescent signal comprise the energy spectrum $S(E) = Gs_0(E)$ (where $s_0$ is normalized to unity, $\int s_0(E) dE = 1$). The emitted energy is isotropic, so that the energy emitted in unit energy interval per unit solid angle is $(G/4\pi) s_0(E)$. The energy $D_i(z)$ reaching the $i$th detector surface ($i=1, 2$) is attenuated in a way that depends on $z$.

In the presence of absorption, characterized by the interband absorption coefficient $\chi_a(E)$, the detection probability for a photon at energy $E$ (averaged over all angles) is given by

$$\pi(E, z) = \int_0^{\infty} \exp[-\chi_a(z) r] \cos \theta \frac{d\theta}{2\pi} dr$$

where $r = z \tan \theta$ and $\theta = z \cos \theta$. Averaged over the emitted photon spectrum $S(E)$, the probability $p_1$ that an emitted photon reaches the 1st photodiode at $z = 0$ is given by

$$p_1(z) = \int \pi(E, z) S(E) dE$$

The probability $(2)$ and the similar quantity for the 2nd photodiode, $p_2(z) = p_1(d-z)$, are referred to as single-pass probabilities, because they do not include the subsequent fate (recycling) of the absorbed photon.

The single-pass probability $p_{\text{fca}}$ of free-carrier absorption can be estimated in a similar fashion, in terms of the FCA absorption coefficient $\chi_a(E)$, viz.

$$p_{\text{fca}} = \int \frac{2\pi}{2\pi + \chi_a} [-\pi(E, z) - \pi(E, d-z)] dE$$

So long as the photon recycling process continues, the minority carriers (holes) and photons are interchangeable entities. The process can be finished off by FCA while the entity is photon or, while the entity is hole, by nonradiative transitions that occur with the small probability $\xi = 1 - \eta$ (in our samples $\eta$ ranges from 90 to 99%).

Fig. 2. Room temperature photon collection efficiency in an InP scintillator of thickness $d = 350 \mu$m, doped to $N_n = 3 \times 10^{17} \text{ cm}^{-3}$, as a function of the event distance $z$ from the surface. The inset shows the ratio of detector signals $D_1/D_2$. The ratio is shown for $z \leq d/2$, since by symmetry, one has $p(D_2) = p(D_1)(d-z)$. 

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site relative to the surface. The problem is how to distinguish the signal arising from a large energy deposited far from the photoreceiver surface from that arising from smaller energy deposited nearby. The problem arises from the attenuation of the optical signal. However, if we knew the distance $z$ of the gamma interaction event from the photoreceiver surface, we could correct for the attenuation.

The solution is based on tallying the signals $D_1$ and $D_2$ individually. The inset to Fig. 2 shows that the ratio of detector signals $p = D_1/D_2 = p_1/p_2$ is an excellent measure of $z$. The simultaneous detection by both detectors of the scintillation arising from the same interaction event, allows us to determine the position of the interaction and therefore correct for attenuation.

3. Conclusion

The efficiency of photon collection in direct-gap semiconductors is limited by parasitic processes, such as FCA and nonradiative recombination of the minority carriers. If these are minimized, one can have an opaque but "ideal" (in terms of the photon collection efficiency) semiconductor scintillator. A nearly ideal scintillator is provided by a lightly-doped InP wafer endowed on both sides with integrated photoreceivers.

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